



If $v_1(p)$ is linear, then

$$\alpha_{eq} = \begin{cases} 1 , v_{1} \leq v_{m} - \Delta v \\ (v_{m} - v_{1}) / \Delta v = (p_{m} - p_{1}) (dv_{1} / dp) / \Delta v , \\ v_{m} - \Delta v \leq v_{1} \leq v_{m} \\ 0 , v_{m} \leq v_{1} \end{cases}$$

Here v_m and p_m are pressure and volume where the Hugoniot first enters the mixed phase. These expressions yield

$$p_{1} = p_{D} + (x \Delta v/2U\tau) dp/dv_{1}, v_{1} \leq v_{m} - \Delta v$$
$$= p_{m} + (p_{D} - p_{m}) \exp(-x/2U\tau),$$
$$v_{m} - \Delta v \leq v_{1} \leq v_{m}$$

 $= 0 \quad v_m \leq v_1$

where x = Ut. Figure (5.17) shows that in the present case (p_D = driving pressure = 200 kbar), the central formula applies, therefore we should expect to find that

$$p_1 - 130 = 70 \exp(-x/2U\tau)$$

The difference between this curve and the numerical results, shown in Fig. 5.13, is due to non-linear effects.

In Figs. 5.5 and 5.11 there are arrows labelled A and B. These indicate the shock front position which would be predicted at the indicated times from the Rankine-Hugoniot jump conditions: A for the first shock, B for the second. The difference between this predicted arrival time and the one obtained in the